Auxiliary Functions Legendre Transforms

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Auxiliary Function

- The term "auxiliary function" usually refers to the functions created during the course of a proof in order to prove the result.
- In thermodynamics, quantities with dimensions of energy were introduced that have useful physical interpretations and simplify calculations in situations where controlled set of variables were used.

Work

- In general, work can be divided into two parts:
 - work of expansion and contraction, and
 - work of the sum of all other forms
- Therefore in the reversible case,

$$\underline{f} \cdot d\underline{X} = -pdV + \sum_{i} \mu_{i} dn_{i}$$

where μ_i will be defined as the chemical potential of species *i*, but not yet at this moment.

Euler's homogeneous function theorem
States that: Suppose that the function *f* is continuously differentiable, then *f* is positive homogeneous of degree n if and only if

$$f(\lambda \underline{x}) = \lambda^n f(\underline{x})$$

 n= 1, f is a first-order homogeneous function

- Let f(x₁,..., x_n) be a first-order homogeneous function of x₁,..., x_n.
- Let $u_i = \lambda x_i$
- Then $f(u_1,...,u_n) = \lambda f(x_1,...,x_n)$
- Differentiate with respect to λ ;

$$\left(\frac{\partial f(u_1,\ldots,u_n)}{\partial \lambda}\right)_{x_i} = f(x_1,\ldots,x_n)$$
(1)

• From calculus,

$$df(u_1, \dots, u_n) = \sum_{i=1}^n \left(\partial f / \partial u_i \right)_{u_j} du_i$$
 (2)

• and,

$$\left(\partial f / \partial \lambda\right)_{x_i} = \sum_{i=1}^n \left(\partial f / \partial u_i\right)_{u_j} \left(\partial u_i / \partial \lambda\right)_{x_i}$$

$$=\sum_{i=1}^{n}\left(\partial f/\partial u_{i}\right)_{u_{j}}x_{i}$$

(3)

• Substitute back to eq. (1),

$$f(x_1, \dots, x_n) = \sum_{i=1}^n \left(\partial f / \partial u_i \right)_{u_j} x_i \qquad (4)$$

• and take $\lambda = 1$,

$$f(x_1, \dots, x_n) = \sum_{i=1}^n \left(\partial f / \partial x_i \right)_{x_i} x_i$$
 (5)

• This is Euler's theorem for first-order homogeneous functions

Legendre Transform

• Recall the 2nd law of thermodynamics,

$$dS = (1/T)dE - (f/T) \bullet dX$$
$$dE = TdS + f \bullet dX$$

- and $\underline{f} \cdot d\underline{X} = -pdV + \sum_{i} \mu_{i}dn_{i}$
- we arrive at, $dE = TdS - pdV + \sum_{i} \mu_{i} dn_{i}$
- Thus, E=E(S,V,n₁,n₂,...n_r), is a natural function of S, V, and the n_i's.

- However, experimentally, *T* is much more convenient than *S*.
- Assume $f = f(x_1, ..., x_n)$ is a natural function of $x_1, ..., x_n$. Euler's theorem
- Then,

Euler's theorem for first-order homogeneous functions

$$f(x_1, \dots, x_n) = \sum_{i=1}^n \left(\partial f / \partial x_i \right)_{x_j} x_i$$

$$df = \sum_{i=1}^{n} u_i dx_i \quad u_i = \left(\frac{\partial f}{\partial x_i}\right)_{x_j}$$

• Let

$$g = f - \sum_{i=r+1}^{n} u_i dx_i$$

• Then,

$$dg = df - \sum_{i=r+1}^{n} (u_i dx_i + x_i du_i)$$
$$= \sum_{i=1}^{r} u_i dx_i + \sum_{i=r+1}^{n} (-x_i) du_i$$

- Thus, $g = g(x_1, ..., x_r, u_{r+1}, ..., u_n)$ is a natural function of $x_1, ..., x_r$ and the conjugate variables to $x_{r+1}, ..., x_n$, namely $u_{r+1}, ..., u_n$.
- The function g is called a Legendre transform of f.

• It transform away the dependence upon $x_{r+1},...,x_n$ to a dependence upon $u_{r+1},...,u_n$.

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• Newton $\vec{F} = m\vec{a}$

• Leibniz
$$\vec{F} = m \frac{d^2 x}{dt^2}$$

• Euler-Lagrange

Hamilton

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial \dot{q}_i} = 0, \quad L = K - V, \text{ 'action'}$$
$$\frac{\partial H}{\partial \dot{q}_i} = -\dot{p}_i$$
$$\frac{\partial H}{\partial \dot{p}_i} = \dot{q}_i, \quad H = K + V, \text{ 'Hamitonian'}$$

See my PPT at https://mestudio.tongji.edu \rightarrow Teaching \rightarrow Grad \rightarrow StatisticalThermo for details

ST.

- It is apparent that this type of transformation allows one to introduce a natural function of *T*, *V*, and *n*, since *T* is simply the conjugate variable to *S*; so as to *p* to *V*.
- From the first and second law, we have E = E(S, V, n)
- We construct a natural function of *T*, *V* and *n*, by subtract from the *E*(*S*, *V*, *n*) the quantity

 $S \times (variable conjugate to S) = ST.$

- Let A(T, V, n) = E TS called the Helmholtz free energy r
- Therefore,

$$dA = -SdT - pdV + \sum_{i=1}^{r} \mu_i dn_i$$

Legendre Transform

- Let G(T, p, n) be the Gibbs free energy G = E - TS - (-pV)
- And H(S, p, n) be the Enthalpy H = E - (-pV) = E + pV
- Therefore,

$$dG = -SdT + Vdp + \sum_{i=1}^{r} \mu_i dn_i$$

Think also about, volume to U pressure to H, natural variables

 $dH = TdS + Vdp + \sum_{i} \mu_{i} dn_{i}$

Maxwell Relations

- Armed with the auxiliary, many types of different measurements can be interrelated.
- Consider,

 $(\partial S / \partial V)_{T,n}$ $\square \square \square$ implies we are viewing S as function of the natural function of T, Vand n.

Maxwell Relations

• If df = adx + bdy, from calculus,

$$\left(\frac{\partial a}{\partial y}\right)_{x} = \left(\frac{\partial b}{\partial x}\right)_{y}$$

- **Recall** $dA = -SdT pdV + \mu dn$
- Then we have

$$(\partial S / \partial V)_{T,n} = (\partial p / \partial T)_{V,n}$$

• and

$$dG = -SdT - Vdp + \mu dn$$
$$\left(\partial S / \partial p\right)_{T,n} = -\left(\partial V / \partial T\right)_{p,n}$$

Example I

• Let $C_v = T(\partial S / \partial T)_{V,n}$

$$dE = TdS - pdV + \sum_{i} \mu_{i} dn_{i}$$

• then $\left(\frac{\partial C_{v}}{\partial V}\right)_{T_{n}} = T\left(\frac{\partial}{\partial V}\left(\frac{\partial S}{\partial T}\right)_{V,n}\right)_{T_{n}}$ $=T\left(\frac{\partial}{\partial T}\left(\frac{\partial S}{\partial V}\right)_{T_n}\right)_{V_n}$ $=T\left(\frac{\partial}{\partial T}\left(\frac{\partial p}{\partial T}\right)_{V,n}\right)_{V,n}$ $=T\left(\frac{\partial^2 p}{\partial T^2}\right)_{-1}$

Quiz X

• Derive an analogous form for (10 Mins)

$$\left(\frac{\partial C_p}{\partial p}\right)_{T,n}$$

- Show that for a one component p-V-n system (10 Mins) $\left(\frac{\partial \mu}{\partial \nu}\right)_T = \nu \left(\frac{\partial p}{\partial \nu}\right)_T$
- where v is the volume per mole. [Hint: show that , where s is the entropy per mole.

$$d\mu = -sdT + vdp$$

Example II

• Let
$$C_p = T\left(\frac{\partial S}{\partial T}\right)_{p,n}$$

- Viewing S as a function of T, V and n
- We have

$$\left(dS \right)_{n} = \left(\frac{\partial S}{\partial T} \right)_{V,n} \left(dT \right)_{n} + \left(\frac{\partial S}{\partial V} \right)_{T,n} \left(dV \right)_{n}$$
$$\left(\frac{\partial S}{\partial T} \right)_{p,n} = \left(\frac{\partial S}{\partial T} \right)_{V,n} + \left(\frac{\partial S}{\partial V} \right)_{T,n} \left(\frac{dV}{\partial T} \right)_{n,p}$$

Maxwell Relations

• Hence
$$\frac{1}{T}C_p = \frac{1}{T}C_v + \left(\frac{\partial p}{\partial T}\right)_{V,n} \left(\frac{\partial V}{\partial T}\right)_{n,p}$$

- Note that $\begin{pmatrix} \frac{\partial x}{\partial y} \\ \frac{\partial y}{\partial z} \end{pmatrix}_{z} = -\left(\frac{\partial x}{\partial z}\right)_{y} \left(\frac{\partial z}{\partial y}\right)_{x} \quad \text{Euler's chain rule}$ So $\left(\frac{\partial p}{\partial T}\right)_{V,n} = -\left(\frac{\partial p}{\partial V}\right)_{T,n} \left(\frac{\partial V}{\partial V}\right)_{p,n}$
 - Therefore

$$C_{p} - C_{v} = -T(\partial p / \partial V)_{T,n} \left[(\partial V / \partial T)_{p,n} \right]$$

• From the 2nd law of thermodynamics,

$$E = E(S, \underline{X})$$

• the internal energy *E* is extensive, it depends upon *S* and <u>*X*</u>, which are also extensive.

$$E(\lambda \underline{X}) = \lambda E(S, \underline{X})$$

• Thus, *E*(*S*,<u>*X*</u>) is a first order homogeneous function of *S* and <u>*X*</u>.

• Therefore, from Euler's theorem, Eq.5,

$$E = \left(\frac{\partial E}{\partial S}\right)_{\underline{X}} S + \left(\frac{\partial E}{\partial \underline{X}}\right)_{\underline{S}} \underline{X}$$
$$= TS + f \cdot \underline{X}$$

where \underline{X} is a vector means system volume

• And work is,

$$\underline{f} \cdot d\underline{X} = -pdV + \sum_{i} \mu_{i} dn_{i}$$

Extensive Function

• This flow naturally as we gave earlier,

$$dE = TdS - pdV + \sum_{i=1}^{r} \mu_i dn_i$$

- That is, $E = E(S, V, n_1, ..., n_r)$
- and Euler's theorem yields,

$$E = TS - pV + \sum_{i=1}^{r} \mu_i n_i$$

Extensive Function

• Its total differential is

$$dE = TdS + SdT - pdV - Vdp + \sum_{i=1}^{r} \left(\mu_i dn_i + n_i d\mu_i \right)$$

• Therefore,

$$0 = SdT - Vdp + \sum_{i=1}^{r} \left(n_i d\mu_i \right)$$

This is the Gibbs-Duhem Equation

Extensive Function

- Recall the definition of Gibbs free energy G = E - TS - (-pV)
- Apply Euler's theorem gives,

$$dG = \left(TS - pV + \sum_{i=1}^{r} \mu_i dn_i\right) - TS - pV$$
$$= \sum_{i=1}^{r} \mu_i dn_i$$

For one component system μ = G/n, Gibbs free energy per mole

Quiz (exercise 1.14)

Show that for a one component p-V-n system

$$\left(\frac{\partial\mu}{\partial\nu}\right)_T = \nu \left(\frac{\partial p}{\partial\nu}\right)_T$$

• where v is the volume per mole. [Hint: show that $d\mu = -sdT + vdp$, where s is the entropy per mole.