

# **Auxiliary Functions**

# **Legendre Transforms**

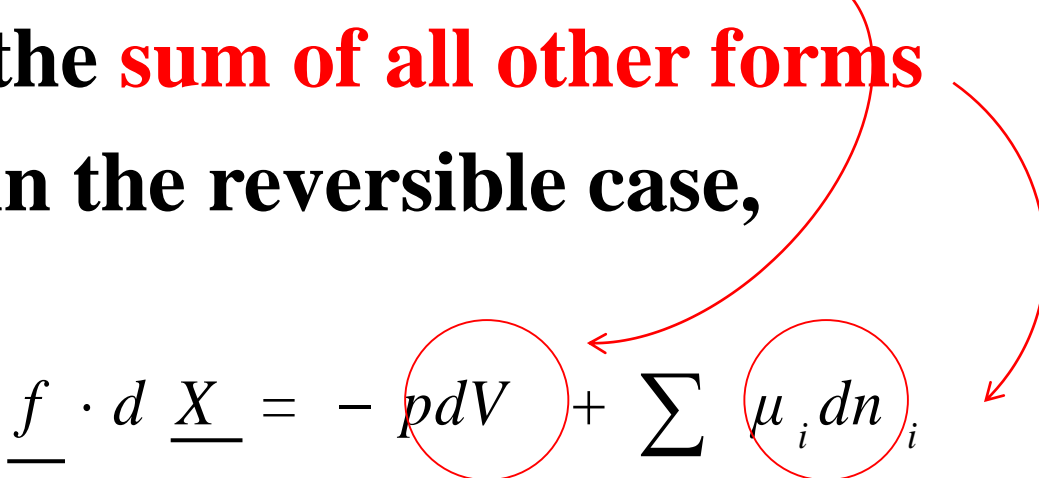
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# Auxiliary Function

- **The term "auxiliary function" usually refers to the functions created during the course of a proof in order to prove the result.**
- **In thermodynamics, quantities with dimensions of energy were introduced that have useful physical interpretations and simplify calculations in situations where controlled set of variables were used.**

# Work

- In general, work can be divided into two parts:
  - work of expansion and contraction, and
  - work of the **sum of all other forms**
- Therefore in the reversible case,

$$\underline{f} \cdot d \underline{X} = - pdV + \sum_i \mu_i dn_i$$


where  $\mu_i$  **will be** defined as the chemical potential of species  $i$ , but **not yet** at this moment.

# Euler's theorem

- **Euler's homogeneous function theorem**

**States that: Suppose that the function  $f$  is continuously differentiable, then  $f$  is positive homogeneous of degree  $n$  if and only if**

$$f(\lambda \underline{x}) = \lambda^n f(\underline{x})$$

- **$n=1$ ,  $f$  is a first-order homogeneous function**

# Euler's theorem

- Let  $f(x_1, \dots, x_n)$  be a first-order homogeneous function of  $x_1, \dots, x_n$ .
- Let  $u_i = \lambda x_i$
- Then  $f(u_1, \dots, u_n) = \lambda f(x_1, \dots, x_n)$
- Differentiate with respect to  $\lambda$ ;

$$\left( \frac{\partial f(u_1, \dots, u_n)}{\partial \lambda} \right)_{x_i} = f(x_1, \dots, x_n) \quad (1)$$

# Euler's theorem

- **From calculus,**

$$df(u_1, \dots, u_n) = \sum_{i=1}^n \left( \frac{\partial f}{\partial u_i} \right)_{u_j} du_i \quad (2)$$

- **and,**

$$\begin{aligned} \left( \frac{\partial f}{\partial \lambda} \right)_{x_i} &= \sum_{i=1}^n \left( \frac{\partial f}{\partial u_i} \right)_{u_j} \left( \frac{\partial u_i}{\partial \lambda} \right)_{x_i} \\ &= \sum_{i=1}^n \left( \frac{\partial f}{\partial u_i} \right)_{u_j} x_i \end{aligned} \quad (3)$$

# Euler's theorem

- **Substitute back to the first equation,**

$$f(x_1, \dots, x_n) = \sum_{i=1}^n \left( \frac{\partial f}{\partial u_i} \right)_{u_j} x_i \quad (4)$$

- **Take  $\lambda = 1$ ,**

$$f(x_1, \dots, x_n) = \sum_{i=1}^n \left( \frac{\partial f}{\partial x_i} \right)_{x_j} x_i \quad (5)$$

- **This is Euler's theorem for first-order homogeneous functions**

# Legendre Transform

- Recall the 2<sup>nd</sup> law of thermodynamics,

$$dS = (1/T)dE - (\underline{f} / T) \cdot d\underline{X}$$

$$dE = TdS + \underline{f} \cdot d\underline{X}$$

- and  $\underline{f} \cdot d\underline{X} = -pdV + \sum_i \mu_i dn_i$
- we arrive at,

$$dE = TdS - pdV + \sum_i \mu_i dn_i$$

- Thus,  $E = E(S, V, n_1, n_2, \dots, n_r)$ , is a **natural function** of **S**, **V**, and the  **$n_i$ 's**.



# Legendre Transform

- **However**, experimentally,  $T$  is much more convenient than  $S$ .

- Assume  $f = f(x_1, \dots, x_n)$  is a natural function of  $x_1, \dots, x_n$ .

- **Then**,  $f(x_1, \dots, x_n) = \sum_{i=1}^n \left( \frac{\partial f}{\partial x_i} \right)_{x_j} x_i$

Euler's theorem for first-order homogeneous functions

$$df = \sum_{i=1}^n u_i dx_i \quad u_i = \left( \frac{\partial f}{\partial x_i} \right)_{x_j}$$

- **Let**  $g = f - \sum_{i=r+1}^n u_i dx_i$

# Legendre Transform

- **Then,** 
$$dg = df - \sum_{i=r+1}^n (u_i dx_i + x_i du_i)$$
$$= \sum_{i=1}^r u_i dx_i + \sum_{i=r+1}^n (-x_i) du_i$$
- **Thus,  $g = g(x_1, \dots, x_r, u_{r+1}, \dots, u_n)$  is a natural function of  $x_1, \dots, x_r$  and the **conjugate variables** to  $x_{r+1}, \dots, x_n$ , namely  $u_{r+1}, \dots, u_n$ .**
- **The function  $g$  is called a Legendre transform of  $f$ .**

# Legendre Transform

- It transform away the dependence upon  $x_{r+1}, \dots, x_n$  to a dependence upon  $u_{r+1}, \dots, u_n$ .
- It is apparent that this type of transformation allows one to introduce a natural function of  $T$ ,  $V$ , and  $n$ , since  $T$  is simply the **conjugate variable** to  $S$ ; so as to  $p$  to  $V$ .

# Legendre Transform

- From the first and second law, we have

$$E = E(S, V, n)$$

- We construct a natural function of  $T$ ,  $V$  and  $n$ , by subtract from the  $E(S, V, n)$  the quantity  $S \times$  (variable conjugate to  $S$ ) =  $ST$ .

- Let  $A(T, V, n) = E - TS$  called the **Helmholtz free energy**

- Therefore,

$$dA = -SdT - pdV + \sum_{i=1}^r \mu_i dn_i$$

# Legendre Transform

- Let  $G(T, p, n)$  be the **Gibbs free energy**

$$G = E - TS - (-pV)$$

- And  $H(S, p, n)$  be **the Enthalpy**

$$H = E - (-pV) = E + pV$$

- Therefore,

$$dG = -SdT + Vdp + \sum_{i=1}^r \mu_i dn_i$$

$$dH = TdS + Vdp + \sum_{i=1}^r \mu_i dn_i$$

Think also about,  
volume to U  
pressure to H

# Maxwell Relations

- Armed with the **auxiliary**, many types of different measurements can be interrelated.
- Consider,

$$\left( \frac{\partial S}{\partial V} \right)_{T, n}$$

-  implies we are viewing  $S$  as function of the natural function of  $T$ ,  $V$  and  $n$ .

# Maxwell Relations

- **If  $df = adx + bdy$ , from calculus,**

$$\left(\frac{\partial a}{\partial y}\right)_x = \left(\frac{\partial b}{\partial x}\right)_y$$

- **Recall**  $dA = -SdT - pdV + \mu dn$

- **Then we have**

$$\left(\frac{\partial S}{\partial V}\right)_{T,n} = \left(\frac{\partial p}{\partial T}\right)_{V,n}$$

- **and**

$$dG = -SdT - Vdp + \mu dn$$

$$\left(\frac{\partial S}{\partial p}\right)_{T,n} = -\left(\frac{\partial V}{\partial T}\right)_{p,n}$$

# Example I

- **Let**  $C_v = T \left( \frac{\partial S}{\partial T} \right)_{V, n}$

- **then** 
$$\left( \frac{\partial C_v}{\partial V} \right)_{T, n} = T \left( \frac{\partial}{\partial V} \left( \frac{\partial S}{\partial T} \right)_{V, n} \right)_{T, n}$$

$$= T \left( \frac{\partial}{\partial T} \left( \frac{\partial S}{\partial V} \right)_{T, n} \right)_{V, n}$$

$$= T \left( \frac{\partial}{\partial T} \left( \frac{\partial p}{\partial T} \right)_{V, n} \right)_{V, n}$$

$$= T \left( \frac{\partial^2 p}{\partial T^2} \right)_{V, n}$$



# Quiz (exercise 1.10)

- **Derive an analogous form for (15 Mins)**

$$\left( \frac{\partial C_p}{\partial V} \right)_{T,n}$$

# solution

$$\begin{aligned} \left( \frac{\partial C_p}{\partial p} \right)_{T, n} &= T \left( \frac{\partial}{\partial p} \left( \frac{\partial S}{\partial T} \right)_{p, n} \right)_{T, n} \\ &= T \left( \frac{\partial}{\partial T} \left( \frac{\partial S}{\partial p} \right)_{T, n} \right)_{p, n} \\ &= T \left( \frac{\partial}{\partial T} \left( - \frac{\partial V}{\partial T} \right)_{p, n} \right)_{p, n} \\ &= -T \left( \frac{\partial^2 V}{\partial T^2} \right)_{p, n} \end{aligned}$$

# Example II

- **Let**  $C_p = T \left( \frac{\partial S}{\partial T} \right)_{p, n}$
- **Viewing  $S$  as a function of  $T$ ,  $V$  and  $n$**
- **We have**

$$(dS)_n = \left( \frac{\partial S}{\partial T} \right)_{V, n} (dT)_n + \left( \frac{\partial S}{\partial V} \right)_{T, n} (dV)_n$$

$$\left( \frac{\partial S}{\partial T} \right)_{p, n} = \left( \frac{\partial S}{\partial T} \right)_{V, n} + \left( \frac{\partial S}{\partial V} \right)_{T, n} \left( \frac{dV}{dT} \right)_{n, p}$$

# Maxwell Relations

- **Hence** 
$$\frac{1}{T} C_p = \frac{1}{T} C_v + \left( \frac{\partial p}{\partial T} \right)_{V, n} \left( \frac{\partial V}{\partial T} \right)_{n, p}$$

- **Note that** 
$$\left( \frac{\partial x}{\partial y} \right)_z = - \left( \frac{\partial x}{\partial z} \right)_y \left( \frac{\partial z}{\partial y} \right)_x$$
 Euler's chain rule

- **So** 
$$\left( \frac{\partial p}{\partial T} \right)_{V, n} = - \left( \frac{\partial p}{\partial V} \right)_{T, n} \left( \frac{\partial V}{\partial T} \right)_{p, n}$$

- **Therefore**

$$C_p - C_v = -T \left( \frac{\partial p}{\partial V} \right)_{T, n} \left[ \left( \frac{\partial V}{\partial T} \right)_{p, n} \right]^2$$

# Euler's theorem

- From the 2<sup>nd</sup> law of thermodynamics,

$$E = E(S, \underline{X})$$

- the internal energy  $E$  is extensive, it depends upon  $S$  and  $\underline{X}$ , which are also extensive.

$$E(\lambda \underline{X}) = \lambda E(S, \underline{X})$$

- Thus,  $E(S, \underline{X})$  is a first order homogeneous function of  $S$  and  $\underline{X}$ .

# Euler's theorem

- Therefore, from Euler's theorem, Eq.5,

$$\begin{aligned} E &= \left( \partial E / \partial S \right)_{\underline{X}} S + \left( \partial E / \partial \underline{X} \right)_s \underline{X} \\ &= TS + \underline{f} \cdot \underline{X} \end{aligned}$$

where  $\underline{X}$  is a vector means system volume

- And work is,

$$\underline{f} \cdot d \underline{X} = -pdV + \sum_i \mu_i dn_i$$

# Extensive Function

- This flow naturally as we gave earlier,

$$dE = TdS - pdV + \sum_{i=1}^r \mu_i dn_i$$

- That is,  $E = E(S, V, n_1, \dots, n_r)$

- and Euler's theorem yields,

$$E = TS - pV + \sum_{i=1}^r \mu_i n_i$$

# Extensive Function

- Its total differential is

$$dE = TdS + SdT - pdV - Vdp + \sum_{i=1}^r (\mu_i dn_i + n_i d\mu_i)$$

- Therefore,

$$0 = SdT - Vdp + \sum_{i=1}^r (n_i d\mu_i)$$

**This is the Gibbs-Duhem Equation**



# Extensive Function

- **Recall the definition of Gibbs free energy**

$$G = E - TS - (-pV)$$

- **Apply Euler's theorem gives,**

$$\begin{aligned} dG &= \left( TS - pV + \sum_{i=1}^r \mu_i dn_i \right) - TS - pV \\ &= \sum_{i=1}^r \mu_i dn_i \end{aligned}$$

- **For one component system  $\mu = G/n$ , Gibbs free energy per mole**

# Quiz (exercise 1.14)

- Show that for a one component p-V-n system

$$\left( \frac{\partial \mu}{\partial v} \right)_T = v \left( \frac{\partial p}{\partial v} \right)_T$$

- where  $v$  is the volume per mole. [Hint: show that  $d\mu = -sdT + vdp$ , where  $s$  is the entropy per mole.]

# solution

- The Gibbs-Duhem Equation,

$$0 = SdT - Vdp + \sum_{i=1}^r (n_i d\mu_i)$$

- Implies, for one component,

$$d\mu = sdT - vdp$$

- Hence,

$$\left( \frac{\partial \mu}{\partial v} \right)_T = -s \left( \frac{\partial T}{\partial v} \right)_T + v \left( \frac{\partial p}{\partial v} \right)_T$$
