Second Law

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Introduction

- Why we interested in the second law ?
- the "arrow of time"
 - It is vividly recognized by consciousness
 - It is equally insisted on by our reasoning faculty (capability to reason)
 - Increase in randomness in the study of organization a number of individuals

Hamilton's Equation

• For a single particle

$$\frac{dp_i}{dt} = -\frac{\partial H}{\partial q_i}$$
$$\frac{dq_i}{dt} = \frac{\partial H}{\partial p_i}$$
$$H = T + V$$
$$= \frac{1}{2m}\mathbf{p}^2 + \mathbf{V}(\mathbf{r})$$
$$H = \sum_{i=1}^3 \frac{p_i^2}{2m} + V(q_1, q_2, q_3)$$

Hamilton's Equations

Energy is a constant of the motion

$$\frac{dH(p_i, q_i)}{dt} = \sum_i \left(\frac{\partial H}{\partial p_i} \frac{dp_i}{dt} + \frac{\partial H}{\partial q_i} \frac{dq_i}{dt} \right)$$
$$\frac{dH}{dt} = \sum_i \left(-\frac{\partial H}{\partial p_i} \frac{\partial H}{\partial q_i} + \frac{\partial H}{\partial q_i} \frac{\partial H}{\partial p_i} \right)$$
$$\frac{dH}{dt} = 0$$

 At microscopic level no sense of "arrow of time", events can equally unfold forward or backward.

- It seams increase in randomness can equally possible forward or backward.
- This lead us to the big band.
- At the beginning of the time, everything is highly ordered, with the expansion, randomness set in.
- Before the "final" state of equilibrium, randomness increases with time.

 With the big ban theory, the Friedmann-Lemaitre-Robertson-Walker (FLRW) metric, we may introduce the Hubble parameter,

$$\mathcal{H} (t) = \frac{d a (t)}{d t}$$
$$a (t)$$

where *a* is a time dependent dimensionless scale factor.

$$\frac{dp_i}{dt} = -\frac{\partial H}{\partial q_i} \qquad H = T + V$$
$$\frac{dq_i}{dt} = \frac{\partial H}{\partial p_i} \qquad = \frac{1}{2m} \mathbf{p}^2 + \mathbf{V}(\mathbf{r})$$

 Using pseudo-Newtonian representation, the new equation of motion is given as

$$\frac{d p}{dt} = -\frac{\partial U(q)}{\partial q} + \frac{d^2 a}{dt^2} \cdot \frac{q}{dt}$$
 force like

• where U(x) is the potential energy

 The i-th component of the observed velocities is then, velocity like d *a* $\frac{dq_{i}}{dt} = \frac{p_{i}}{m_{i}} + \frac{dt}{a} \cdot q_{i} = \frac{p_{i}}{m_{i}} + \mathcal{H}$

· q_i

- Obviously the last term breaks the time symmetry, when $(t \rightarrow -t)$, \prec $\frac{dq_{i}}{dt} = \frac{p_{i}}{m_{i}} - \mathcal{H} \cdot q_{i}$
- metric expansion of space gives a small correction to the Hamiltonian time evolution of the system.

 For times much shorter than 1/λ, (1/λ is the Lyapunov time, the characteristic timescale on which a dynamical system is chaotic), this small correction to the Hamiltonian time evolution result in the entropy difference is,

$$\Delta S \approx k_{B}(t - t_{0}) \int \rho(\underline{x}_{0}, \underline{p}_{0}) \Lambda d \underline{x}_{0} d \underline{p}_{0}$$

- Then $\Delta S \approx k_B (t t_0) NH_0$
- This difference is always positive. This means that the direction of the deviation from the Hamiltonian path due to FLRW expansion is always towards an entropy increase.

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 Postulate: There is an extensive function of state, S(E,X), which is a monotonically increasing function of E, and if state B is adiabatically accessible from state A, then

$$S_{B} \geq S_{A}$$

- Notice that if this state B was reversibly accessible from state A, this postulate also implies that $S_A = S_B$
- If A and B are adiabatically and reversibly accessible,

$$S_A \ge S_B$$

The extensive function of state, S(E, X) is called the entropy.

 $dS = (\partial S / \partial E)_X dE + (\partial S / \partial X)_E \cdot dX$

 $dE = (dQ_{rev} + f \cdot dX)$

 $dS = (\partial S / \partial E)_{\underline{X}} (dQ_{rev} + \left[(\partial S / \partial \underline{X})_E + (\partial S / \partial E)_{\underline{X}} f \right] \cdot d\underline{X}$

• For an adiabatic process that is reversible, both dS and $(dQ)_{rev}$ are zero

 Since, all displacements connecting the manifold of equilibrium states

 $(\partial S / \partial \underline{X})_E = -(\partial S / \partial E)_{\underline{X}} \underline{f}$

 Notice all quantities involved in this equation are functions of state, therefore, it holds for nonadiabatic as well as adiabatic processes.

 In the postulate, it states that S is a monotonically increasing function of E;

 $(\partial S/\partial E)_X > 0, \text{ or } (\partial E/\partial S)_X \ge 0$

Let $T \equiv (\partial E / \partial S)_X \ge 0$

 Note that both *E* and *S* are extensive, the temperature is intensive, that is, independent of the size of the system.

The definition of temperature

• Substitute back,

$$(\partial S / \partial X)_E = -f / T$$

• Therefore,

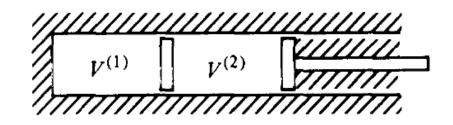
$$dS = (1/T)dE - (f/T) \cdot dX$$

$$dE = TdS + f \cdot dX$$

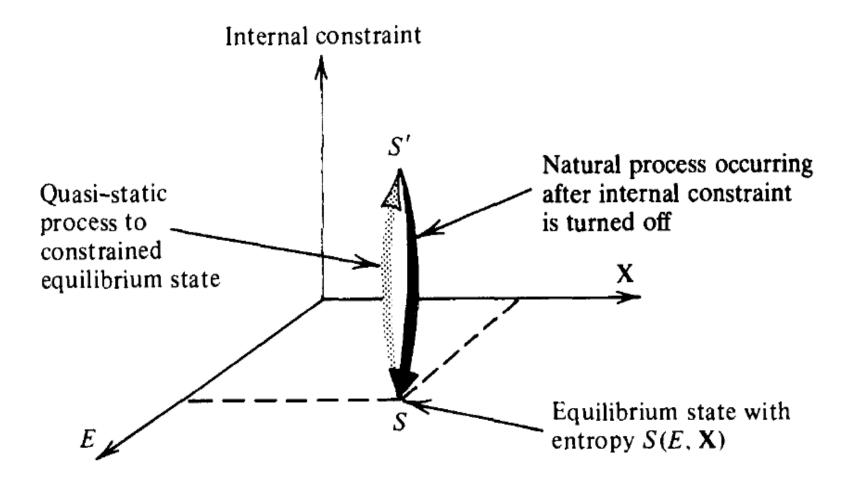
$$(\Delta S)_{adiabatic} \ge 0$$

 These equations constitute the mathematical statement of the second law

 Internal constraints: are constraints that couple to extensive variables but not alter the total value of those extensive variables (e.g. partition)



- The system is initially at equilibrium S = S(E, X)
- Then, by applying an internal constraint, the system is reversibly brought to a constrained equilibrium with the same E and <u>X</u>, but with entropy S' = S(E, <u>X</u>; internal constraint)

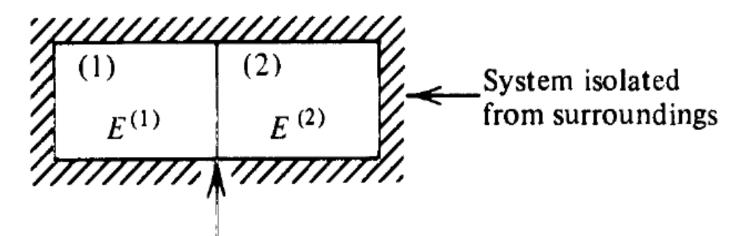


- The states on the *E*-*X* plane are the manifold of equilibrium states in the absence of the internal constraint
- The application of the internal constraint lifts the system off this manifold. It will require work to do this
- and the requirement that there is no change in energy for the process means that there must also have been a flow of heat

- After attaining and while maintaining this constrained state, the system will be adiabatically insulated.
- Then, the internal constraint is suddenly shut off
- The system will relax naturally at constant *E* and <u>X</u> back to the initial state with entropy S

- According to the second law, the entropy change is positive S - S' > 0,
- Or S(E, X) > S(E, X; internal constraint)—the equilibrium state is the state at which S(E, X; internal constraint) has its global maximum

The Energy minimum principle:



Heat conducting wall divides subsystems (1) and (2)

Repartitioning the system

The entropy Maximum requires that

 $S(E^{(1)} - \Delta E, \underline{X}^{(1)}) + S(E^{(2)} + \Delta E, \underline{X}^{(2)}) < S(E^{(1)} + E^{(2)}, \underline{X}^{(1)} + \underline{X}^{(2)})$ after repartition before

- Entropy is extensive, △E is an amount of energy removed from subsystem 1 and placed into subsystem 2.
- Since S is a monotonically increasing function of *E*, (increases in E will led to increase in S, and vice versa, monotonic)
- Therefore

• When $\Delta E \neq 0$, and allow

$$S(E^{(1)} - \Delta E, \underline{X}^{(1)}) + S(E^{(2)} + \Delta E, \underline{X}^{(2)}) = S(E, \underline{X}^{(1)} + \underline{X}^{(2)})$$

- To raise left hand from less to equal, energy must be added
- That is, apply internal constraints at constant S and <u>X</u>, will necessarily raise the total energy of the system
- Therefore, the equilibrium state *E*(*S*, <u>*X*</u>) is the state at which *E*(*S*, <u>*X*</u>; internal constraint) has its global minimum

Scenario I: at equilibrium

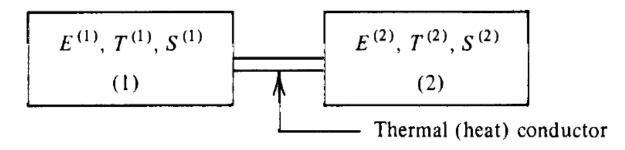


Fig. 1.6. Heat conducting system.

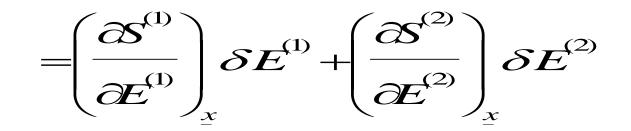
• How are T⁽¹⁾ and T⁽²⁾ related?

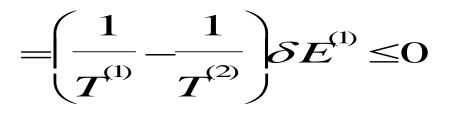
 Consider a small displacement about equilibrium due to an internal constraint,

$$(\delta S)_{E,\underline{X}} \leq 0$$
 (entropy maximum)

- Since $E = E^{(1)} + E^{(2)}$ is constant during the displacement (see diagram), $\delta E^{(1)} = -\delta E^{(2)}$
- Since S is extensive, $S = S^{(1)} + S^{(2)}$

• Thus $\delta S = \delta S^{(1)} + \delta S^{(2)}$





For all δE , therefore, $T^{(1)} = T^{(2)}$

- Scenario II: initially not at equilibrium, then eventually reach equilibrium
- From the second law,

$$\Delta S^{(1)} + \Delta S^{(2)} = \Delta S > 0$$

Assuming differences are small

$$\left(\frac{\partial S^{(1)}}{\partial E^{(1)}}\right)_{X} \Delta E^{(1)} + \left(\frac{\partial S^{(2)}}{\partial E^{(2)}}\right)_{X} \Delta E^{(2)} > 0$$

• Then

$$\Delta S = \left(\frac{1}{T^{(1)}} - \frac{1}{T^{(2)}}\right) \Delta E^{(1)} > 0$$

- Therefore,
 - $T^{(1)} > T^{(2)} \rightarrow \Delta E^{(1)} < \mathbf{0}$
 - $T^{(1)} < T^{(2)} \rightarrow \Delta E^{(1)} > \mathbf{0}$
 - That is energy flow is from the hot body to cold body